

**ANSWERS TO ASSIGNED EXERCISES 1-59**

1. (d.)

$$\lambda = \frac{1}{2}B_0xy$$

2.

$$\begin{aligned}4\pi\epsilon_0\vec{E} &= \frac{q}{r^2}\hat{r} \\ \rho &= q\delta^3(\vec{r}) \\ \vec{B} &= 0 \\ \vec{J} &= 0\end{aligned}$$

3.

$$\begin{aligned}4\pi\epsilon_0V' &= \frac{q}{r} \\ \vec{A}' &= 0\end{aligned}$$

5.

It is always possible to find a gauge in which  $V' = 0$ , by using  $\lambda = \int_0^t V(t') dt'$ . But, in general, a gauge cannot be found in which  $\vec{A}' = 0$ , because that would force  $\vec{B} = 0$ .

6.

- (b.) The constant of proportionality is  $1/3!$ .  
(c.)

$$\det A = \frac{1}{4!}\epsilon_{ijkl}A_{im}A_{jn}A_{kp}A_{lq}\epsilon_{mnpq}$$

10.

- (a.) No (the interval is spacelike).  
(b.) Yes.

$$|\vec{r}_A'' - \vec{r}_B''| = \sqrt{2}$$

- (c.) No (the interval is spacelike).  
(d.) No (the interval is timelike).  
(e.) Yes.

$$c|t_E'' - t_D''| = \sqrt{7}$$

11.  $\Lambda$  is symmetric, and its independent elements are

$$\begin{pmatrix} \gamma & -\gamma\beta n_1 & -\gamma\beta n_2 & 0 \\ \gamma n_1^2 + n_2^2 & (\gamma - 1)n_1n_2 & 0 & 0 \\ & \gamma n_2^2 + n_1^2 & 0 & 0 \\ & & & 1 \end{pmatrix}$$

12. The answers are the same.

15.

$$\begin{aligned}\eta_{\max} &= 10.34 \\ \beta_{\max} &= 1 - (2.09 \times 10^{-9}) \\ x_{\max} &= 2.84 \times 10^{20} \text{ m} = 29,900 \text{ light yr} \\ \Delta t &= 1.89 \times 10^{12} \text{ sec} = 59,850 \text{ yr}\end{aligned}$$

16. (b.)

$$\text{fraction decaying} = 1 - \exp(-L/\gamma_0\beta_0c\tau)$$

17.

(a.)

$$L = \gamma\beta c\tau = \frac{4}{3}c\tau$$

(b.)

$$\psi = \cos^{-1}(2\beta^2 - 1) = 73.74^\circ$$

19. (b.) The dimensions of  $d\sigma$  are Joules<sup>-2</sup>.

24.

(a.)

$$I = 2ne\beta_0cA$$

(b.)

$$\frac{1}{\mu_0}\vec{B} = \hat{\phi} \frac{ne\beta_0cA}{\pi s}$$

(c.)

$$n'_+ = n/\gamma_0$$

(d.)

$$n'_- = n\gamma_0(1 + \beta_0^2)$$

(e.)

$$c\epsilon_0\vec{E}' = -\hat{s}\gamma_0\beta_0\frac{ne\beta_0cA}{\pi s}$$

26.

(a.)

$$F_{\mu\nu}F^{\mu\nu} = -\frac{2}{c^2}(|\vec{E}|^2 - |c\vec{B}|^2)$$

(b.)

$$F^{\mu\nu}G_{\mu\nu} = -\frac{4}{c}\vec{E} \cdot \vec{B}$$

(c.)

$$\vec{E} \perp \vec{B} \quad \text{and} \quad |\vec{E}| > |c\vec{B}|$$

**28.**

(a.)

$$\sinh \eta = \frac{F_0}{mc} t$$

(b.)

$$t_1 = \frac{mc}{F_0}$$

**29.**(c.)  $x$  does not increase linearly with  $t$ .**30.**(c.)  $z$  does increase linearly with  $t$ .**32.**

$$\vec{E}(\vec{r}, t) = \hat{z} \frac{\mu_0 q_0 c^3 t}{2\pi[(ct)^2 - s^2]^{3/2}} \quad \text{for } t > s/c$$

$$= 0 \quad \text{otherwise}$$

$$c\vec{B}(\vec{r}, t) = -\hat{\phi} \frac{\mu_0 q_0 c^2 s}{2\pi[(ct)^2 - s^2]^{3/2}} \quad \text{for } t > s/c$$

$$= 0 \quad \text{otherwise}$$

**35.**

(b.)

$$\vec{S} \rightarrow \frac{q^2 c}{16\pi^2 \epsilon_0} \frac{\hat{z}}{b^2 z^2} \quad \text{as } \gamma \rightarrow \infty$$

(c.) No.

**37.**

(a.)

$$4\pi\epsilon_0 V(\vec{r}) \cong \frac{4ak}{\pi r}$$

(b.)

$$4\pi\epsilon_0 V(\vec{r}) \cong \frac{2a^2 k \cos \theta}{\pi r^2}$$

(c.)

$$4\pi\epsilon_0 V(\vec{r}) \cong -\frac{2a^3 k}{\pi^2 r^3} (3 \cos^2 \theta - 1)$$

**38.**

(a.)

$$\frac{\epsilon_0}{q} V(\vec{r}) = \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{Y_{lm}(\theta, \phi) Y_{lm}^*(\theta', \phi') r'^l}{(2l+1)r^{l+1}}$$

(b.)

$$\frac{4\pi\epsilon_0}{q} V(\vec{r}) \cong \frac{1}{r}$$

$$+ \frac{r'}{r^2} (\cos \theta \cos \theta' + \sin \theta \sin \theta' \cos(\phi - \phi'))$$

$$+ \frac{r'^2}{r^3} \left( \left( \frac{3}{2} \cos^2 \theta - \frac{1}{2} \right) \left( \frac{3}{2} \cos^2 \theta' - \frac{1}{2} \right) \right.$$

$$+ 3 \sin \theta \sin \theta' \cos \theta \cos \theta' \cos(\phi - \phi')$$

$$\left. + \frac{3}{4} \sin^2 \theta \sin^2 \theta' \cos(2(\phi - \phi')) \right)$$

**39.**Place charges  $+1, -4, +6, -4, +1$  at  $z = 2, 1, 0, -1$ , and  $-2$ , respectively.**42.**

$$f(\theta, \phi) \propto \cos^2 \theta \sin^2 \theta$$

**43.**(a.)  $q_{22}$  and  $q_{2-2}$  do not vanish and have equal relative weight.(b.) E-type (TM) radiation of types  $E_{22}$  and  $E_{2-2}$  are emitted with equal relative weight.

(c.)

$$f(\theta, \phi) \propto \sin^2 \theta (1 - \sin^2 \theta \cos^2 2\phi)$$

vanishes in six directions:  $\theta = 0, \theta = \pi$ , and at  $(\theta = \frac{\pi}{2}; \phi = 0, \frac{\pi}{2}, \pi, \frac{3\pi}{2})$ .**47.**

(a.)

$$\omega' = \frac{2\pi\beta c}{\Delta z/\gamma} \quad (\gamma \equiv \frac{1}{\sqrt{1-\beta^2}})$$

(b.)

$$\omega = \frac{\omega'}{\gamma(1-\beta)}$$

(c.)

$$\begin{aligned}\lambda &= \Delta z \frac{1 - \beta}{\beta} \\ &= \frac{\Delta z}{\gamma^2 \beta (1 + \beta)} \\ &\rightarrow \frac{\Delta z}{2\gamma^2} \text{ as } \beta \rightarrow 1\end{aligned}$$

(d.)

$$\begin{aligned}E &= \gamma m_e c^2 \approx \sqrt{\frac{\Delta z}{2\lambda}} m_e c^2 \\ &= 1.616 \text{ GeV}\end{aligned}$$

**51.**

(b.) The state of polarization is RH (LH) elliptical if  $\text{Im } \beta < 0$  ( $\text{Im } \beta > 0$ ).

(c.)

$$\begin{aligned}\sqrt{1 + |\beta|^2} \vec{J}_1 &= \begin{pmatrix} 1 \\ \beta \end{pmatrix} \\ &= -\text{Im } \beta \begin{pmatrix} 1 \\ -i \end{pmatrix} + \begin{pmatrix} 1 + \text{Im } \beta \\ \text{Re } \beta \end{pmatrix} \\ &= \text{Im } \beta \begin{pmatrix} 1 \\ i \end{pmatrix} + \begin{pmatrix} 1 - \text{Im } \beta \\ \text{Re } \beta \end{pmatrix}\end{aligned}$$

No, this concern would not invalidate the answer to (b.).

**52.**

$$\begin{aligned}I_{A+B} &= I_A + I_B + \\ &+ \sqrt{I_A I_B} \frac{\text{Re}(\alpha \gamma^* + \beta \delta^*)}{\sqrt{(|\alpha|^2 + |\beta|^2)(|\gamma|^2 + |\delta|^2)}}\end{aligned}$$

**53.**

(a.)

$$J = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

(c.) No, the twisted nematic cell absorbs  $\hat{y}$  polarized light, whereas the rotator rotates  $\hat{y}$  polarized light into  $\hat{x}$  polarized light.

**56.**

(a.) No, Pedrotti's Jones matrix completely absorbs LH circularly polarized light.

(b.) Light passes through a device consisting of

a QWP with slow axis at  $45^\circ$  with respect to  $\hat{x}$ , followed by an  $\hat{x}$  polarizer, followed by a QWP with slow axis at  $-45^\circ$ . The Jones matrix for this combination is half of that in part (a.).

**59.**

(a.)

$$\begin{aligned}\begin{pmatrix} \mathcal{S}_0 \\ \mathcal{S}_1 \\ \mathcal{S}_2 \\ \mathcal{S}_3 \end{pmatrix} &\equiv \mathcal{S}_0 = \mathcal{S}_p + \mathcal{S}_n \\ &= \begin{pmatrix} \sqrt{\mathcal{S}_1^2 + \mathcal{S}_2^2 + \mathcal{S}_3^2} \\ \mathcal{S}_1 \\ \mathcal{S}_2 \\ \mathcal{S}_3 \end{pmatrix} + \\ &+ \begin{pmatrix} \mathcal{S}_0 - \sqrt{\mathcal{S}_1^2 + \mathcal{S}_2^2 + \mathcal{S}_3^2} \\ 0 \\ 0 \\ 0 \end{pmatrix}\end{aligned}$$